

Study of anomalous gauge couplings using $W^\pm Z$ channel in ATLAS

Study of dim-6 effective field theory in vector boson scattering

Lilly Wüst

March 25, 2019



Searching for new physics at the LHC

- since there is nothing new found in direct searches
 - indirect searches become important
 - searching for deviations of the SM and parametrizing them with an effective field theory

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \sum_{i,j} \left[\frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \frac{f_j}{\Lambda^4} \mathcal{O}_j^{(8)} \right] \quad \text{with } \Lambda \text{ scale of new physics}$$

Effective field theory

Expand the standard model Lagrangian with operators of higher dimension:

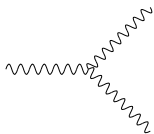
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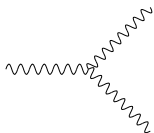


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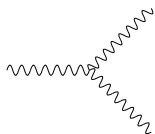
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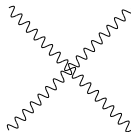
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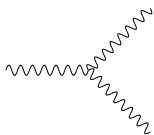


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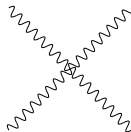
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dim-8 operators:



- \rightarrow introduce additional interactions between the gauge bosons
- \rightarrow anomalous triple and quartic gauge couplings
- \rightarrow will modify the cross section and kinematic distribution

Current status of dim-6 limits

- LEP, Tevatron, ATLAS and CMS already provided limits on dim-6 operators
- as a reference I will use 13 TeV ATLAS limits with 13.3 fb^{-1} [arXiv:1603.02151v1]:

EFT coupling	Observed [TeV^{-2}]
c_{www}/Λ^2	$[-3.6; 3.4]$
c_w/Λ^2	$[-4.1; 7.6]$
c_b/Λ^2	$[-261.0; 193.0]$

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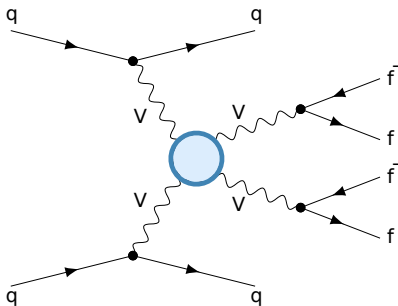
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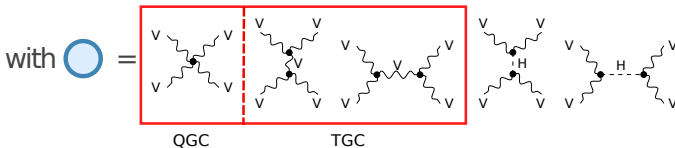
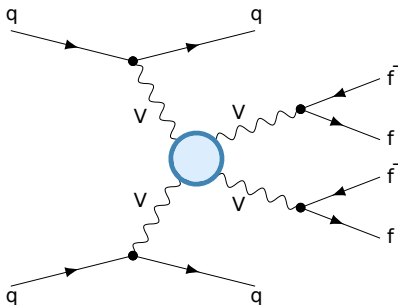
- as comparison 13 TeV CMS limits with 35.9 fb^{-1} [arXiv:1901.03428v1]:

EFT coupling	Observed [TeV^{-2}]
c_{www}/Λ^2	$[-2.0; 2.1]$
c_w/Λ^2	$[-4.1; 1.1]$
c_b/Λ^2	$[-100.0; 160.0]$

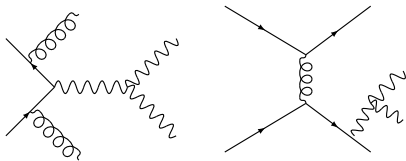
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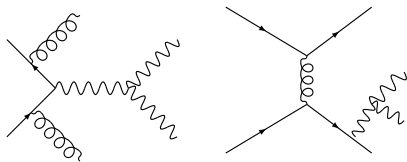


The $VVjj$ production:



and a lot more

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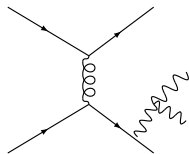
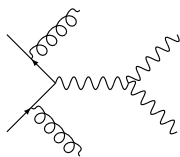


and a lot more

$$\propto \alpha_w^4 \alpha_s^2$$

→ *QCD* mediated production

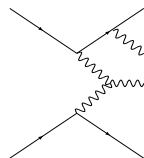
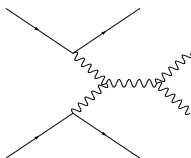
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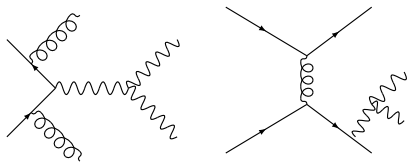
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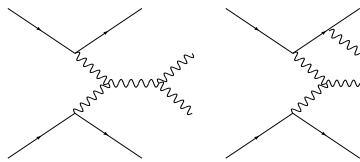
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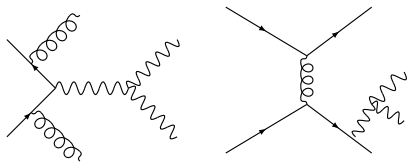


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→ EW mediated production

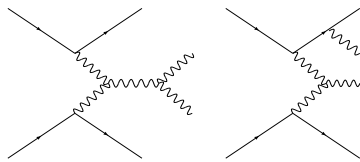
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and a lot more

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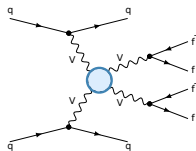
$$\propto \alpha_w^6$$

→ EW mediated production

→ classify $VVjj$ production in two kinds

The VBS phase space

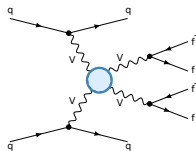
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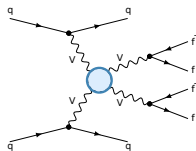
- at least two jets
- the two jets with highest p_T and opposite hemisphere are the tagging jets
- no b-jets
- $M_{inv}^{jj} > 500 \text{ GeV}$



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→ reduce many QCD and EW mediated production of $VVjj$

Motivation of my study

- ideally calculate limits on dim-6 and dim-8 operators together
 - strong dim-6 limits LEP and LHC already produced using inclusive $W^\pm Z$
 - to set limits on the dim-8 operators VBS was typically used at the LHC
 - assuming that dim-6 operators have a negligible impact
- I want to study the impact of dim-6 in the VBS phase space for dim-8 operators

Simulation of $W^\pm Z$ production

I simulated $W^\pm Z$ in the fully leptonic decay channel with madgraph on truth level for the SM and dim-6 operators:

- for $W^\pm Z$ with *QCD* mediated production:

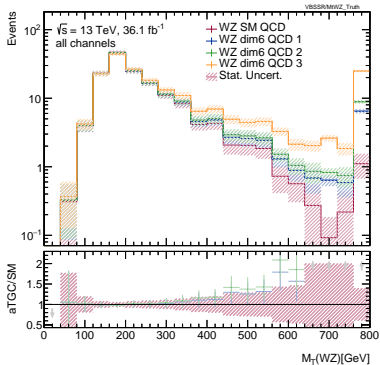
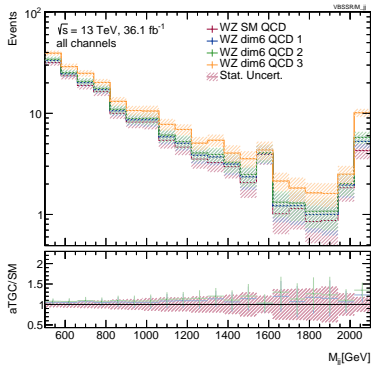
$$pp \rightarrow W^\pm Z \rightarrow \ell\ell\nu$$

with up to two jets in the matrix element

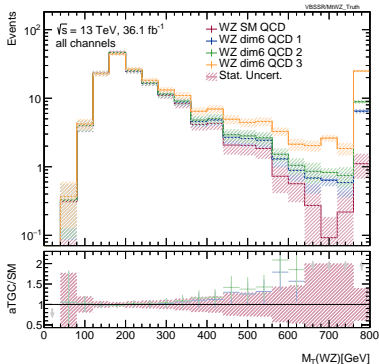
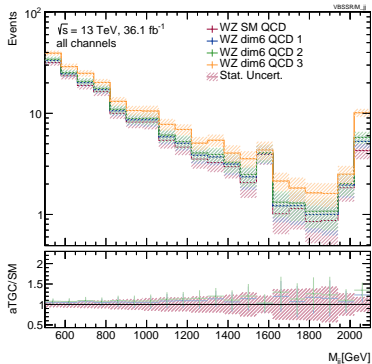
- for $W^\pm Z$ with *EW* mediated production:

$$pp \rightarrow W^\pm Z \rightarrow \ell\ell\nu jj$$

The Distributions



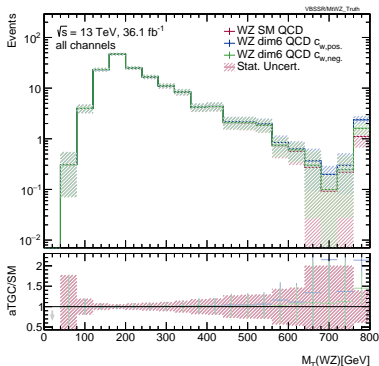
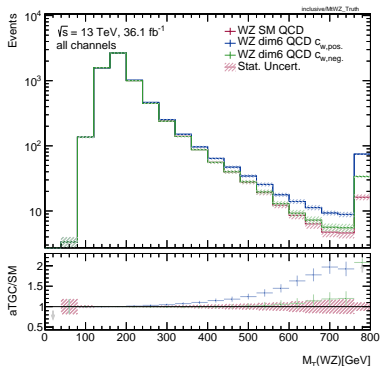
The Distributions



since M_T^{WZ} is the most sensitive distribution I will use it

M_T^{WZ} in the limits

M_T^{WZ} distribution in the limit of c_W of aTGC for QCD mediated production of $W^\pm Z$:



→ calculate the sum of number events and take the ratio $N_{aTGC}^{QCD} / N_{SM}^{QCD}$ to get an idea of the dim-6 operators impact



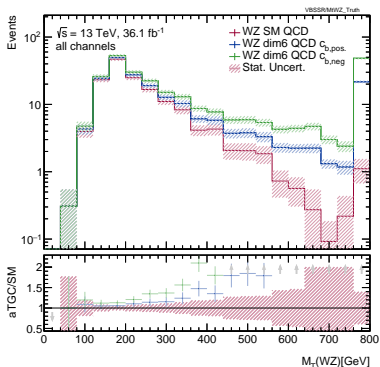
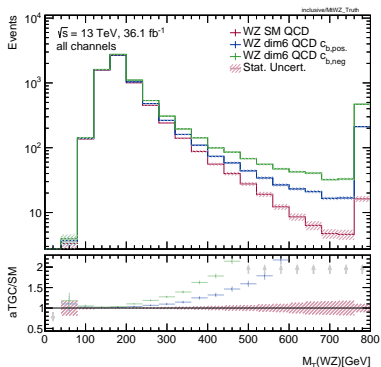
M_T^{WZ} in the limits

- without VBS cuts: $N_{SM}^{QCD} = 6497$ events
- with VBS cuts: $N_{SM}^{QCD} = 153$ events

	without VBS cuts	with VBS cuts
	$N_{aTGC}^{QCD} / N_{SM}^{QCD}$	$N_{aTGC}^{QCD} / N_{SM}^{QCD}$
$C_{www,neg.}$	1.011	1.022
$C_{www,pos.}$	1.013	1.028
$C_{W,neg.}$	1.003	1.011
$C_{W,pos.}$	1.026	1.012
$C_{b,neg.}$	1.383	1.739
$C_{b,pos.}$	1.232	1.322

M_T^{WZ} in the c_b limits

M_T^{WZ} distribution in the limit of c_b of aTGC for QCD mediated production of $W^\pm Z$:



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- have a look on the impact of dim-8 operators in the VBS phase space
- compare the dim-6 and dim-8 operators impact
- calculate the uncertainty of the dim-8 limits and the cross section σ with the impact of dim-6 operators

Thank you for your attention!

Dim-6 operators

X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	Q_φ	$(\varphi^\dagger \varphi)^3$	$Q_{e\varphi}$	$(\varphi^\dagger \varphi)(\bar{l}_p e_r \varphi)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$Q_{\varphi\Box}$	$(\varphi^\dagger \varphi)\Box(\varphi^\dagger \varphi)$	$Q_{u\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p u_r \tilde{\varphi})$
Q_W	$\varepsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$Q_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D_\mu \varphi)$	$Q_{d\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p d_r \varphi)$
$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{l}_p \gamma^\mu l_r)$
$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{l}_p \tau^I \gamma^\mu l_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}_p \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}_p \gamma^\mu u_r)$
$Q_{\varphi WB}$	$\varphi^\dagger \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{d}_p \gamma^\mu d_r)$
$Q_{\varphi \tilde{W}B}$	$\varphi^\dagger \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{u}_p \gamma^\mu d_r)$

Dim-6 operators

$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
Qu	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	Q_{ie}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{dd}	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	Q_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		B -violating			
Q_{ledq}	$(\bar{l}_p^j e_r)(\bar{d}_s^j q_t^k)$	Q_{duq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^{\gamma j})^T C l_t^k]$		
$Q_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	Q_{qqq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	Q_{qqq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jkmn} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^{\gamma m})^T C l_t^n]$		
$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	Q_{duu}	$\varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$				



The inclusive $W^\pm Z$ phase space:

- less than four leptons passing the baseline selection*
- three leptons passing the Z lepton selection*
- leading lepton $p_T > 27 \text{ GeV}$
- two leptons with same flavour and opposite charge passing the Z lepton selection*
- mass window: $|M_{\ell\ell} - M_Z| < 10 \text{ GeV}$
- remaining lepton passing W lepton selection*
- W-lepton and p_T^{miss} have $m_T^W > 30 \text{ GeV}$

* [ATLAS-Note STDM-2018-03]