

Coupling Deformation and Flow in Fractured Poroelastic Media

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Fluid–solid coupling

- \blacktriangleright Porous material, contains thin layer with different material properties.
- Coupled processes:
	- 1. Fluid-fluid: The fluid diffuses from the fracture into the porous medium and vice versa.
	- 2. Fluid-solid: The pressure applied by the fluid onto the fracture faces induces a deformation.
	- 3. Solid-fluid: The flow in the fracture is affected by the fracture aperture *b* and thus on the deformation of the surrounding solid skeleton.
	- 4. Poroelasticity: The matrix fluid pressure interacts with matrix stress
- **►** Generally $b \ll a$ \Rightarrow use dimension-reduced flow model.

Darcy law in $\Omega \setminus \overline{\Sigma}$

$$
\nabla\cdot\mathbf{q}^\Omega=f^\Omega
$$

$$
-\mathbf{K}\nabla\rho^\Omega=\mathbf{q}^\Omega
$$

$$
\nabla \cdot \mathbf{q}^{\Omega} = f^{\Omega}
$$

$$
-\mathbf{K} \nabla \rho^{\Omega} = \mathbf{q}^{\Omega}
$$

Darcy law in $\Omega \setminus \overline{\Sigma}$ Averaged Darcy law on Σ [Martin, Jaffré, Roberts '05]

$$
\nabla_{\tau} \cdot \mathbf{q}^{\Sigma} - [\![\mathbf{q}^{\Omega}]\!] \cdot v = bt^{\Sigma}
$$

$$
-b\mathbf{K}^{\tau} \cdot \nabla_{\tau} p^{\Sigma} = \mathbf{q}^{\Sigma}
$$

Coupling constraints on Σ , $\xi \in (1/2, 1]$

$$
\begin{aligned} \left\{ \mathbf{q}^{\Omega} \right\} \cdot \mathbf{v} &= \frac{-K^{\mathbf{v}}}{b} \left[p^{\Omega} \right] \\ \left[\mathbf{q}^{\Omega} \right] \cdot \mathbf{v} &= \frac{4}{2\xi - 1} \frac{K^{\mathbf{v}}}{b} \left(p^{\Sigma} - \left\{ p^{\Omega} \right\} \right) \end{aligned}
$$

$$
\nabla \cdot \mathbf{q}^{\Omega} = f^{\Omega}
$$

$$
-\mathbf{K} \nabla \rho^{\Omega} = \mathbf{q}^{\Omega}
$$

Darcy law in $\Omega \setminus \overline{\Sigma}$ Averaged Darcy law on Σ [Martin, Jaffré, Roberts '05]

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$$

$$
-b\mathbf{K}^{\tau} \cdot \nabla_{\tau} p^{\Sigma} = \mathbf{q}^{\Sigma}
$$

$$
\begin{aligned} \mathbf{0} &= \nabla \cdot \left(\sigma(\mathbf{u}^{\Omega}) - \rho^{\Omega} \mathbb{I} \right) \\ \epsilon(\mathbf{u}^{\Omega}) &= \frac{1}{2} \left(D \mathbf{u}^{\Omega} + (D \mathbf{u}^{\Omega})^{\mathsf{T}} \right) \\ \sigma(\mathbf{u}^{\Omega}) &= \mathbb{E} : \epsilon(\mathbf{u}^{\Omega}) \end{aligned}
$$

Poroelasticity equations in $\Omega \setminus \overline{\Sigma}$ Coupling constraints on Σ , $\xi \in (1/2,1]$

$$
\begin{aligned} \left\{ \mathbf{q}^{\Omega} \right\} \cdot \mathbf{v} &= \frac{-K^{\nu}}{b} \left[p^{\Omega} \right] \\ \left[\mathbf{q}^{\Omega} \right] \cdot \mathbf{v} &= \frac{4}{2\xi - 1} \frac{K^{\nu}}{b} \left(p^{\Sigma} - \left\{ p^{\Omega} \right\} \right) \\ - p^{\Sigma} \mathbf{v}^{\pm} &= \sigma \left(\mathbf{u}^{\Omega} \right) \cdot \mathbf{v}^{\pm} \\ b &= \left[\mathbf{u}^{\Omega} \right] \cdot \mathbf{v} \end{aligned}
$$

Poroelastic problem

For fixed fluid pressures p^{Ω} , p^{Σ} :

 \blacktriangleright Symmetric bilinear form:

$$
a_E^{\Omega}(\mathbf{u}^{\Omega},\mathbf{v}^{\Omega})=\int_{\Omega}\sigma(\mathbf{u}^{\Omega}):\varepsilon(\mathbf{v}^{\Omega}).
$$

Coupling terms:

$$
c_{E}^{\Omega}(\rho^{\Omega}; \mathbf{v}^{\Omega}) = -\int_{\Omega} \nabla \rho^{\Omega} \cdot \mathbf{v}^{\Omega} \cdot \mathbf{v} d\Gamma
$$

$$
c_{E}^{\Sigma}(\rho^{\Sigma}; \mathbf{v}^{\Omega}) = \int_{\Sigma} 2\rho^{\Sigma} \mathbf{v}^{\Omega} \cdot \mathbf{v} d\Gamma
$$

Define

$$
V:=\left\{\nu\in L^2(\widetilde{\Omega})\,\bigg|\, \nu|_{\Omega^\pm}\in H^1(\Omega), [\![\nu]\!]_{\Sigma_\varrho}=0\right\}.
$$

and $V_E \subset V^d$ as the space of functions satisfying the Dirichlet boundary conditions

► Weak formulation: Find **such that**

$$
a_{E}^{\Omega} (\textbf{u}^{\Omega},\textbf{v}^{\Omega}) + c_{E}^{\Omega} (\rho^{\Omega};\textbf{v}^{\Omega}) + c_{E}^{\Sigma} (\rho^{\Sigma};\textbf{v}^{\Omega}) = \textbf{0} \hspace{1cm} \forall \textbf{v}^{\Omega} \in \textbf{V}_{0}
$$

 \blacktriangleright Linear elasticity problem with pre-stress

For fixed fluid pressures p^{Ω} , p^{Σ} :

Theorem ([Folklore])

 \mathcal{A} *ssume* $\nabla p^{\Omega} \in L^2(\Omega)$, $p^{\Sigma} \in H^{\frac{1}{2}}(\Sigma)$ *and* $\textbf{u}_D^{\Omega} \in \textbf{H}^{\frac{1}{2}}(\Gamma_D^E)$. Then there exists a unique *Solution* $u^{\Omega} \in V_E$ *b to the linear elasticity problem with pre-stress. Furthermore the solution can be decomposed additively into a continuous function u^c and a singular part us. Near the crack tip the singular part can be approximated* b *v* $u_s = K_l u_l + K_l u_l + K_l u_l$ *u_{II}*, where

$$
u_{I,i}, u_{II,i}, u_{III,i} \in span\left\{\sqrt{r}\sin\frac{\Theta}{2}, \sqrt{r}\cos\frac{\Theta}{2}, \sqrt{r}\sin\Theta\sin\frac{\Theta}{2}, \sqrt{r}\sin\Theta\cos\frac{\Theta}{2}\right\},\right\}
$$

for i = 1,...,*d where* (*r*,Θ) *are the crack tip polar coordinates.*

Fluid–fluid problem

For fixed fracture width *b*:

 \blacktriangleright Symmetric, bilinear forms:

$$
a_F^{\Omega}(b; p^{\Omega}, r^{\Omega}) = \int_{\Omega} \mathbf{K} \nabla p^{\Omega} \cdot \nabla r^{\Omega} d\Omega + \int_{\Sigma} \frac{K^{V}}{b} [\rho^{\Omega}] [\Gamma^{\Omega}] d\Gamma
$$

$$
+ \frac{4}{2\xi - 1} \int_{\Sigma} \frac{K^{V}}{b} {\rho^{\Omega}} {\{\Gamma^{\Omega}\}} d\Gamma,
$$

$$
a_F^{\Sigma}(b; p^{\Sigma}, r^{\Sigma}) = \int_{\Sigma} b \mathbf{K}^{\tau} \nabla_{\tau} p^{\Sigma} \cdot \nabla_{\tau} r^{\Sigma} + \frac{4}{2\xi - 1} \frac{K^{V}}{b} p^{\Sigma} r^{\Sigma} d\Gamma,
$$

 \blacktriangleright Linear forms:

$$
I_F^{\Omega}(r^{\Omega}) = \int_{\Omega} f_F^{\Omega} r^{\Omega} d\Omega \quad \text{and} \quad I_F^{\Sigma}(b; r^{\Sigma}) = \int_{\Sigma} b f_F^{\Sigma} r^{\Sigma} d\Gamma.
$$

 \blacktriangleright Coupling terms:

$$
c_F(b; r^{\Omega}, \rho^{\Sigma}) = -\frac{4}{2\xi - 1} \int_{\Sigma} \frac{K^V}{b} \rho^{\Sigma} \left\{ r^{\Omega} \right\} d\Gamma,
$$

Coupled weak problem - For fixed fracture width *b*: Find $(\rho^{\Omega}, \rho^{\Sigma})$, such that

$$
a_F^{\Omega}(b; p^{\Omega}, r^{\Omega}) + c_F(b; r^{\Omega}, p^{\Sigma}) = l_F^{\Omega}(r^{\Omega}) \quad \forall r^{\Omega} \in H_0^1(\Omega)
$$

$$
a_F^{\Sigma}(b; p^{\Sigma}, r^{\Sigma}) + c_F(b; p^{\Omega}, r^{\Sigma}) = l_F^{\Sigma}(r^{\Sigma}) \quad \forall r^{\Sigma} \in H_0^1(\Sigma)
$$

► Linear (if *b* is fixed!)

Existence and uniqueness?

- **►** Simple but restrictive result: Existence and uniqueness for $0 < c \leq b(x) \leq C$ a.e. on Σ
- ► Crack front in domain \implies *b*(*x*) \rightarrow 0 for *x* approaching the crack front

Existence and uniqueness with crack tips:

- ► Crack tip asymptotics yield: $b(x) = \text{dist}(x, \gamma)^{\frac{1}{2}}$ near the crack tip and bounded away from zero otherwise
- **►** Denote by $L^2_{b^{-1}}(Σ)$ and $L^2_b(Σ)$ the sets of all measurable functions *v* : $Ω → ℝ$ for which the norms

$$
||v||_{0,b^{-1},\Sigma}^2 := \int_{\Sigma} |v|^2 b^{-1} dx < \infty,
$$

$$
||v||_{0,b,\Sigma}^2 := \int_{\Sigma} |v|^2 b dx < \infty,
$$

respectively. The spaces $\left(L_{b^{\pm 1}}^2(\Sigma),\|\cdot\|_{0,\pm,\Sigma} \right)$ are Hilbert spaces

We have

$$
L^2_{b^{-1}}(\Sigma)\hookrightarrow L^2(\Sigma)\hookrightarrow L^2_b(\Sigma)
$$

Crack-weighted Sobolev spaces:

 \blacktriangleright The spaces

$$
H_b^1(\Sigma) := \left\{ s \in L_{b^{-1}}^2(\Sigma) \mid \nabla_{\tau} s \in L_b^2(\Sigma) \right\},
$$

$$
H_{b^{-1}}^{\frac{1}{2}}(\Sigma) := \left\{ s \in H^{\frac{1}{2}}(\Sigma) \mid s \in L_{b^{-1}}^2(\Sigma) \right\},
$$

$$
H_b^1(\Omega) := \left\{ v \in V \mid \gamma^{\pm} v \in H_{b^{-1}}^{\frac{1}{2}}(\Sigma) \right\},
$$

together with the norms

$$
||s||_{1,b,\Sigma}^2 := ||s||_{0,b^{-1},\Sigma}^2 + ||\nabla_{\tau}s||_{0,b,\Sigma}^2,
$$

$$
||s||_{\frac{1}{2},b^{-1},\Sigma}^2 := ||s||_{0,b^{-1},\Sigma}^2 + \int_{\Sigma \times \Sigma} \frac{|s(x) - s(y)|^2}{|x - y|^d} dx dy,
$$

$$
||v||_{1,b^{-1},\Omega}^2 := ||\gamma^+ v||_{0,b^{-1},\Sigma}^2 + ||\gamma^- v||_{0,b^{-1},\Sigma}^2 + ||u||_1^2
$$

are Hilbert spaces.

 V Define V _{*F*} := V ^E_{*F*} × V ^Ω_{*F*}</sub> (where V ^E_{*F*} ⊂ *H*₁¹_{*b*}</sub>(Σ) and V ^Ω_{*F*} ⊂ *H*₁¹_{*b*}(Ω)) as the spaces of functions satisfying the Dirichlet boundary data.

For fixed crack width *b*:

Theorem ([Hanowski '16])

Assume that Σ *is a bounded parametrized surface with smooth boundary and that the permeability tensor K is symmetric, bounded and uniform elliptic. Let b* > 0 *on* Σ *a.e. and* $b \simeq$ *dist*(\cdot , γ)^{$\frac{1}{2}$} *near the crack tip and* $\xi \in (\frac{1}{2}, 1]$. Let $\Gamma_D^F \neq \emptyset$, $f_F^{\Omega} \in L^2(\Omega)$, $f_F^{\Sigma} \in L^2_b(\Sigma)$, $q_N^{\Omega} \in L^2(\Gamma_N)$ and $q_N^{\Sigma} \in L^2_b(\Sigma)$. Furthermore, assume that

$$
\rho_D^{\Omega} \in W_D^{\Omega} := \{ s \in H^{\frac{1}{2}}(\Gamma_D^F) \mid E_{\Omega} s \in V_b \}
$$

and

$$
\rho_D^\Sigma\in W_D^\Sigma:=\{s\in H^\frac{1}{2}(\gamma_D^F)\ |\ E_\Sigma s\in H_b^1(\Sigma)\},
$$

 ω *where* $E_{\Omega}: H^{\frac{1}{2}}(\Gamma_D^F) \to V$ and $E_{\Sigma}: H^{\frac{1}{2}}(\gamma_D^F) \to H^1(\Sigma)$ are the standard extension *operators. Then there exists a unique solution* $(p^{\Omega}, p^{\Sigma}) \in V_F^{\Omega} \times V_F^{\Sigma}$ *of the weak coupled fluid–fluid problem.*

Fully Coupled Problem

Fixed-point formulation:

 \triangleright Solution operator for the fluid problem :

$$
S_f: H^{\frac{1}{2}}_{00}(\Sigma) \to V_F^{\Omega} \times V_F^{\Sigma}, \quad b \mapsto (p^{\Omega}, p^{\Sigma})
$$

 \triangleright Solution operator for the elasticity problem:

$$
S_e: V_F^\Omega \times V_F^\Sigma \to V_E, \quad (\rho^\Omega, \rho^\Sigma) \mapsto u^\Omega
$$

 \triangleright Normal jump operator

$$
j: \mathbf{V}_{E} \to H_{00}^{1/2}(\Sigma), \quad \mathbf{u}^{\Omega} \mapsto b := [\![\mathbf{u}]\!] \cdot \mathbf{v}
$$

 \blacktriangleright Fixed Point Formulation:

$$
b=(j\circ S_e\circ S_f)b
$$

Existence and Uniqueness:

► Dependency of the solution spaces for the pressures on the fracture width *b*! Iterative approach leads to 'independent' solution spaces in each iteration step.

Challenge:

- **►** Pressure *p*^Ω and displacement **u**^Ω discontinuous on Σ
- ► Derivatives of *p*^Ω and **u**^Ω singular at crack front

XFEM basic idea:

- \blacktriangleright Additional bespoke FE functions near the crack
- \blacktriangleright Reproduce discontinuity: Heavyside function
- \blacktriangleright Reproduce singularity: Special singularity functions derived by asymptotic analysis
- \blacktriangleright Aim: Retain optimal discretization errors

Given:

- \triangleright Grids for $\overline{\Omega}$ and $\overline{\Sigma}$ and/or implicit crack representation
- \blacktriangleright Element subsets:
	- \blacktriangleright \mathcal{K}_1 : elements fully cut by the fracture
	- \triangleright \mathcal{K}_2 : elements in the vicinity of the crack tip

Crack Front Enrichment Functions

Crack front enrichment functions:

- \blacktriangleright (r, Θ) front polar coordinates
- \triangleright Standard enrichment functions for the displacement

$$
(F_E)_1(r, \Theta) = \sqrt{r} \sin\left(\frac{\Theta}{2}\right)
$$

$$
(F_E)_2(r, \Theta) = \sqrt{r} \cos\left(\frac{\Theta}{2}\right)
$$

$$
(F_E)_3(r, \Theta) = \sqrt{r} \sin\left(\frac{\Theta}{2}\right) \sin(\Theta)
$$

$$
(F_E)_4(r, \Theta) = \sqrt{r} \cos\left(\frac{\Theta}{2}\right) \sin(\Theta)
$$

 \blacktriangleright Laplace enrichment for the matrix pressure:

$$
(F_F)_1(r,\Theta) = \sqrt{r}\sin\left(\frac{\Theta}{2}\right)
$$

$$
(F_F)_2(r,\Theta) = \sqrt{r}\cos\left(\frac{\Theta}{2}\right)
$$

Algebraic problem:

- \blacktriangleright Linear coupled fluid–fluid problem
- \blacktriangleright Linear elasticity problem with pre-stress
- ▶ Nonlinear coupling through changing crack width *b*

Decouple by fixed-point iteration

Iteration variable: fracture width function b_k

$$
\begin{array}{ll}\n\text{Init: } b_0; \\
\text{while } (\textit{err} > \textit{acc}) \text{ do} \\
& \begin{cases} (p_{k+1}^{\Omega}, p_{k+1}^{\Sigma}) \leftarrow \text{Solve fluid problem with } b_{\mathbf{u}} = b_k; \\
\mathbf{u}_{k+1}^{\Omega} \leftarrow \text{Solve the elasticity problem with } p^{\Omega} = p_{k+1}^{\Omega} \text{ and } p^{\Sigma} = p_{k+1}^{\Sigma} \text{;} \\
& b_{k+1} \leftarrow b_k + \beta \left([\![\mathbf{u}_{k+1}^{\Omega}]\!] \cdot \mathbf{v} - b_k \right); \\
\text{end}\n\end{cases}
$$

- Damping parameter $\beta \in (0,1]$
- Solve two linear, symmetric systems in each iteration step
- Direct subdomain solvers

Grid

Instructured triangle grid for the bulk, uniform grid for the fracture

Boundary Conditions

► Solid: Zero Dirichlet and Neumann boundary conditions; Fluid: $\rho_0^{\Sigma} = 0.5$ MP, zero Dirichlet and Neumann boundary conditions otherwise

Material Properties:

- ► Bulk Domain: $Ω = [0, 1] \times [-\frac{1}{2}, \frac{1}{2}]$ km²
- ► Fracture: $\Sigma = [0, \frac{1}{2}] \times \{0\}$ km
- \blacktriangleright Fluid: Homogeneous and isotropic permeability tensors, with $K = 0.1$ mD for the bulk and $K^{\rm V} = K^{\tau} = 100{\rm D}$
- \triangleright Solid: St. Venant–Kirchhoff material law with Young's modulus $E = 1$ GP and Poisson ratio $v = 0.3$

Parameters

- **Damping parameter:** $\beta = 1$
- Discretization parameter $\xi = 0.75$

Measuring the Solver Convergence:

 \triangleright Constant rates, fast convergence, nearly independent of the mesh size!

Discretization Error Rates

Measuring the Discretization Error

- \triangleright Compare with solution on fine grid
- \triangleright Optimal error rates!

3d matrix + 2d fracture:

- \blacktriangleright Unit cube
- \blacktriangleright Half-penny crack